The interpretation of very low frequency Raman scattering from roughened silver electrodes: Clusters, cavities, or complexes?

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The potential dependence of the broad low frequency (5–8 cm⁻¹) Raman scattering features from roughened silver electrodes and their sensitivity to the chemical composition of the supporting electrolyte are presented. We find that the electromagnetic enhancement theories which relate these Raman bands to the vibrations of the silver surface structures are unable to account for all of the data; electronic coupling to the stabilizing adsorbates requires that some of the charge transfer concepts of the adatom complex theories be included in the model.

I. INTRODUCTION

Surface roughness has been recognized from the first to be a key factor in the light scattering process from electrochemical systems.¹ An oxidation–reduction cycle (ORC) is required in all but a few instances to see “surface enhancement effects” in surface enhanced Raman scattering (SERS),² second harmonic generation,³ and Rayleigh scattering⁴ from electrodes. A crucial question which immediately arises when discussing surface roughness is the length scale involved in the surface enhancement. Two roughness scales have been proposed to explain the surface enhancement: atomic scale and large scale or geometric (> 50 nm) roughness. In the latter, classical electromagnetic equations show how surface roughness allows the incoming light to couple to the localized surface plasmon modes of the metal, giving rise to an enhancement in the available electromagnetic field amplitude at the surface on the order of 100.⁵

Experiments with attenuated total reflection (ATR),⁶ gratings,⁷ colloids,⁸ and microlithographic structures⁹ verify that at least part of the enhancement of light scattering from surfaces with predesigned topologies is due to these large scale features. By implication it is assumed that similar effects operate at the surface of electrodes roughened by an ORC.

There are, however, several observations for electrochemical systems which the electric field enhancement mechanisms fail to explain. In particular, the SERS signal is lost irreversibly when the electrode is cathodically polarized.¹⁰ Moreover, the intensities of various Raman bands can vary significantly as a function of potential and laser wavelength.¹¹ To explain these effects a second, atomic scale roughness has been proposed.² Adatom complexes on the surface which strongly couple to the molecular adsorbates electronically (and therefore chemically) are postulated to be involved in the SERS active site.² This adatom model has the implicit assumption that the SERS activity is due to a small subpopulation of molecules on the surface. This assumption has been supported by the destruction of the SERS signal by the deposition of a submonolayer (< 3%) of lead or thallium atoms on the silver surface.¹²,¹³ The fact that only a small percentage of the surface molecules contribute to the Raman signal implies that perhaps both the geometric and atomic roughness enhancements are necessary to observe SERS signals.

The recently proposed “cavity site” model¹⁴,¹⁵ presents the possibility of combining both of these effects into a single picture of the surface morphology. In the thermal desorption studies in UHV evidence is cited which suggests that some of the pyridine molecules adsorbed on coldly deposited silver films are transported in low temperature annealing processes to the cracks or cavities which occur at the grain boundaries or other defect-like regions. These cavity sites could show enhanced Raman signals from both electromagnetic effects and strong chemical bonding due to the large anisotropy of the cavity size. It is conceivable that similar surface structures could exist in the aqueous electrochemical systems as a result of the activating ORC.¹⁵

In electrochemical systems the best evidence for some type of surface species or cluster has been obtained by Weitz et al.¹⁶ They observed a broad and inhomogeneous Raman band at very low frequency which they labeled an “inelastic Mie scattering” mode, attributing this band to the antisymmetric breathing vibration of large silver clusters formed at the surface of the electrode during anodization. The observed frequencies (5–10 cm⁻¹) imply a silver cluster size of about 10 nm in diameter, which is on the low end of the geometric roughness scale mentioned previously. An electromagnetic field enhancement for the clusters was cast in the form of a coincidental matching of the laser frequency and the localized plasmon mode, and this enhancement mechanism was used to explain dependence of the bond position on the laser excitation frequency. In addition to the inelastic Mie scattering, Weitz et al.¹⁶ also reported an additional contribution to the low frequency Raman band which they labeled “inelastic Rayleigh scattering” from the molecules coupled to these low frequency cluster vibrations. This last component was found to be very sensitive to the chemical composition of the electrolyte, and was particularly large in the presence of pyridine.

We have critically examined the sensitivity of these low frequency Raman bands (which for lack of a better name we shall refer to as cluster bands) to the potential of the silver electrode and to the chemical composition of the electrolyte. The cluster bands have a potential dependence similar (but not identical) to the molecular SERS bands which have been observed. The Raman signal is irreversibly lost as the potential is swept negatively towards the point of hydrogen evolution, but there is also a quasireversible increase in the band intensity with the same potential change. In addition, we
observe an increase in the signal if a small amount of hydrogen is evolved from the electrode prior to the ORC. This fact and the observation that the cluster bands are critically sensitive to the electrolyte employed—but only when that electrolyte is present during the ORC—lead us to conclude that the cluster bands of these roughened electrodes are due to some unusual surface structure which is stabilized by the adsorbrates and ligands from the solution. This strong interaction between ligand and silver atom implies that the notions from the adatom model of charge transfer and electronic reorganization of these large metal–ligand complexes must be included in our picture of these surface species.

II. EXPERIMENTAL CONSIDERATIONS

All solutions were prepared from deionized water, reagent grade potassium bromide or sodium chloride, and spectroscopic grade pyridine. The solutions were deaerated with Ar prior to electrochemical use. A PAR 173 potentiostat and PAR 175 programmer were used to control a standard three electrode potentiostat which had a Pt wire counter electrode and a Ag/AgCl reference electrode; all potentials stated here are vs Ag/AgCl.

The working electrode in the potentiostat was roughened solid silver electrode which was polished with 0.5 μ alumina, rinsed with deionized water, and inserted into the electrochemical system. The electrode was then roughened by an ORC consisting of a potential step from −400 to +200 mV and back again. The time spent at +200 mV was adjusted so that a charge density of 25 mC/cm² was delivered to the electrode.

The beam from a Coherent model 3000 K krypton ion laser was focused onto the roughened electrode with an incident angle of 45 deg, and the Raman scattering was collected normal the silver surface. The Raman spectra were obtained by a Spex 14018 double monochromator in tandem with a third grating stage (also supplied by Spex) which provided the necessary stray light rejection. The resolution was set at 0.1 cm⁻¹ for recording spectra, and was opened to a 0.5 cm⁻¹ bandwidth for the potential dependence studies.

III. RESULTS AND DISCUSSION

A. Cathodic polarization of the electrode

Polarization of a roughened silver electrode to a sufficiently negative potential irreversibly destroys the molecular SERS signal. However, a large enhancement of the signal is observed in the same systems if hydrogen is evolved at the electrode prior to an ORC. We observe the same effects in the cluster bands. Curve (a) in Fig. 1 shows the low frequency Raman spectrum of a silver electrode in a 0.1 M KBr + 0.02 M pyridine solution (an electrolyte comparable to those employed in the previous study) which has been roughened with a mild ORC (25 mC/cm²) after 10 s at −1400 eV. An additional cathodic polarization for 5 s at the same potential destroyed the signal irreversibly as witnessed in a subsequent spectrum [curve (c) of Fig. 1]. The ability to enhance or destroy the Raman signal reaffirms the belief that the bands are due to some volatile surface species, whether it be a metal cluster, cavity site, or adatom complex.

![Figure 1. Low frequency inelastic light scattering from a roughened silver electrode in a 0.1 M KBr + 0.02 M pyridine electrolyte. Incident laser power: 20 mW; resolution: 0.1 cm⁻¹; electrode potential: −400 mV vs Ag/AgCl. Curve (a): Stokes and anti-Stokes Raman scattering after the electrode was polarized to −1400 mV for 10 s and then subjected to an ORC which delivered 25 mC/cm² to the electrode; laser wavelength: 530.9 nm. Curve (b): same as curve (a) except with a laser wavelength of 647.1 nm. Curve (c): same as curve (a) except that the electrode potential was held at −1400 mV for 5 s after the ORC.](http://jcp.aip.org/jcp/copyright.jsp)
other ORC. These results are in sharp contrast to the molecular SERS spectra. In the bromide/chloride system, the Ag–Cl stretching band at 240 cm$^{-1}$ is replaced by a Ag–Br stretching band at 190 cm$^{-1}$ upon the addition of bromide to a 0.1 M chloride electrolyte. This displacement agrees with differential capacitance studies which show that bromide is much more strongly adsorbed to a silver electrode than chloride ions.

The inability to chemically alter the Raman signal suggests that the chemical enhancement of the signal comes primarily from the particular molecule's (or ion's) role in the production/stabilization of the surface species. The fact that the signal does not change also means that the ligands involved in the surface species are behaving differently than the other ligands adsorbed onto the electrode, i.e., there is strong chemical interaction and bonding occurring at the surface structure ligand sites. In terms of the adatom complex model this interaction allows for possibility of charge transfer processes and resonant effects.

As a final note we mention in passing that laser irradiation effects during the ORC appear to make a small contribution to the enhancement of the production of surface structures, but they were not examined in detail here.

C. The band shape of the cluster band

One of the most intriguing observations made by Weitz et al. was that the position of the Raman cluster band shifted with the laser excitation wavelength. This shift is seen in curve (b) of Fig. 1, which was taken with the 647.1 nm line of the Kr$^+$ laser. Weitz et al. attributed the shift to the selection (by the laser) of the localized plasmon frequency of a different subpopulation of metal clusters present on the electrode (presumably similar plasmon matching conditions could be obtained from the surface cavity sites). Their interpretation implies that a change in band shape with potential, if observed, would mean the selective dissolution of certain metal clusters. We found no change in the band shape as a function of potential for the systems we examined. However, Genack et al. in an earlier paper, did see a potential dependent band shape from a solution of pyrazine.

An alternative explanation for the shift in frequency is to allow for a large distribution of charge transfer complexes which one tunes through as the laser frequency is changed. The Raman signal is then due to this distribution of silver cluster types, each subpopulation having its own charge transfer bands and electronic-vibrational coupling terms. Such a picture would also give a large inhomogeneous band shape that depended on the laser wavelength.

D. The potential dependence of the cluster band

Although we found no change in the band shape with potential for a given laser excitation wavelength, we did observe changes in the overall intensity of the cluster band. However, care must be taken when measuring the amount of signal as a function of potential since the irreversible changes in the cluster band's intensity also affect the spectra. This is seen in Fig. 2. The spectrum at the bottom of the figure was taken at an electrode potential of $-400$ mV immediately following an ORC. The next spectrum, at $-700$ mV, shows an overall increase in the band intensity. However, a subsequent spectrum at $-550$ mV has only roughly half the intensity of the original spectrum at $-400$ mV; this inability to recover the total band intensity that was initially present is due to the irreversible loss of signal while sitting at $-700$ mV. Further spectra at $-850$ mV, and again at $-700$ mV after another ORC, illustrate how the signal intensity depends crucially on the electrode's prior history. Since some band intensity is lost as a function of time at every potential, the notion of a "reversible" potential dependence of the band intensity is meaningless.

We can, however, by carefully choosing our experiment, show that there is an additional change in intensity with potential that occurs on a faster timescale than the irreversible signal loss. Figure 3 plots the signal observed when we set the spectrometer to the peak of the cluster band (with a bandwidth of 0.5 cm$^{-1}$) and then cycle the electrode potential between $400$ and $800$ mV. We see that each time the potential is swept toward $800$ mV the band intensity increases, but that due to the loss of signal as a function of potential there is a gradual loss of band intensity. This short time, or "quasireversible" behavior, occurs at both 530.9 and 647.1 nm laser excitation wavelengths. It can be most
this signal change as due to the shifts in the various charge transfer levels of the complex.

For comparison with the cluster band's potential dependence, we have measured by an identical experiment the potential dependence of the molecular pyridine band at 1008 cm$^{-1}$ (with a spectrometer bandwidth of 5 cm$^{-1}$). The results, plotted in Fig. 4, show striking differences in both the quasi reversible and irreversible behavior between the two molecular Raman spectra taken with 530.9 and 647.1 nm laser excitation wavelengths. These differences have been noted previously by other authors.\(^2\)

IV. CONCLUSION

From our experiments we can make some delineating statements as to the nature of the cluster band. First, the presence of a broad, inhomogeneous band at very low frequencies suggests that there is a distribution of surface structures, involving many silver atoms and ligands, which is responsible for the signal. Whether these structures should be thought of as spheroidal clusters, cavities or dendrites is still subject to debate. It is clear, however, that these structures are not responsible for all of the SERS signal from these roughened electrodes. The disappearance of the Ag–Cl stretch and the subsequent appearance of the Ag–Br stretch in the absence of any change in the cluster band signal leads us to conclude that the existence of SERS active sites is not intimately tied to the existence of these surface species, which are specific surface structures created during the ORC. The effect of molecules and/or ions from solution on the cluster band intensity is crucial if not completely dominant. However, these chemical effects are mainly the stabilization of the clusters during the ORC. The surface species giving rise to the signal in these electrochemical systems must be thought of containing significant amount of interaction between the silver atoms and the adsorbed ligands. This statement is supported by the quasi reversible changes with potential of the intensities of the various molecular and cluster SERS bands. Therefore, the concepts of the adatom model must be incorporated into our picture of these surface species to completely explain the observed volatility of the Raman signal.

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